

# Influence de la nature de la paroi sur les réactions de recombinaison des radicaux

A. Granier, M. Carette, M. Richard, A. Goullet

A. Bousquet (PhD 2005), S. Elisabeth (PhD 2015), Master students

Luc Stafford

Université   
de Montréal

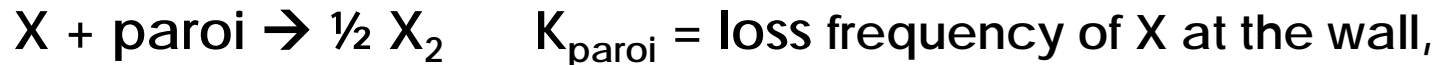
Institut des Matériaux Jean Rouxel, CNRS Université de Nantes, France



# OUTLINE

- ✦ Introduction
- ✦ Kinetics of atoms in low pressure molecular plasmas
- ✦ How to determine the atom recombination frequency ( $\gamma$ ) ?
- ✦ Example : determination of  $\gamma_{\text{O}}$  by time resolved optical emission spectroscopy (TROES)
  - influence of the nature of the wall on  $\gamma_{\text{O}}$  in pulsed oxygen ICP plasmas
  - Values of  $\gamma_{\text{O}}$  in  $\text{O}_2/\text{TTIP}$  and  $\text{O}_2/\text{HMDSO}$  plasmas
- ✦ Conclusion

# Kinetics of X atoms in low pressure plasmas



$$K_{\text{paroi}} = 1/\tau, \quad \tau : \text{characteristic loss time}$$

## Global Model :

$$[X] = 2 K_{\text{diss}} \times n_e \times [X_2] / K_{\text{paroi}} = \tau \times K_{\text{diss}} \times n_e \times [X_2]$$

**$X_2$  dissociation degree :**

$$= \frac{[X]}{2[X_2] + [X]} = \frac{100}{1 + \frac{2K_{\text{paroi}}}{K_{\text{diss}} n_e}}$$

# Recombination probability : $\gamma$

$$\tau = \tau_{diff} + \tau_{loss}$$

$$\tau = \frac{\Lambda^2}{D} + \frac{V}{A} \cdot \frac{2(2-\gamma)}{\bar{v} \cdot \gamma}$$

$$\frac{1}{\Lambda^2} = \left(\frac{\pi}{L}\right)^2 + \left(\frac{2.405}{R}\right)^2$$

$D$  = binary diffusion coefficient of X in X<sub>2</sub>

$V$  = Volume of the reactor

$A$  = Area of the reactor

$\bar{v}$  = thermal velocity of X

Chantry et al J. Appl. Phys. 62, 1141 (1987)

# Dissociation degree of $\text{Cl}_2$ as a function of $\gamma_{\text{Cl}}$

V.M. Donnelly, J. Guha, L. Stafford, JVSTA. 010801(2011)  $n_e$  (cm<sup>-3</sup>)

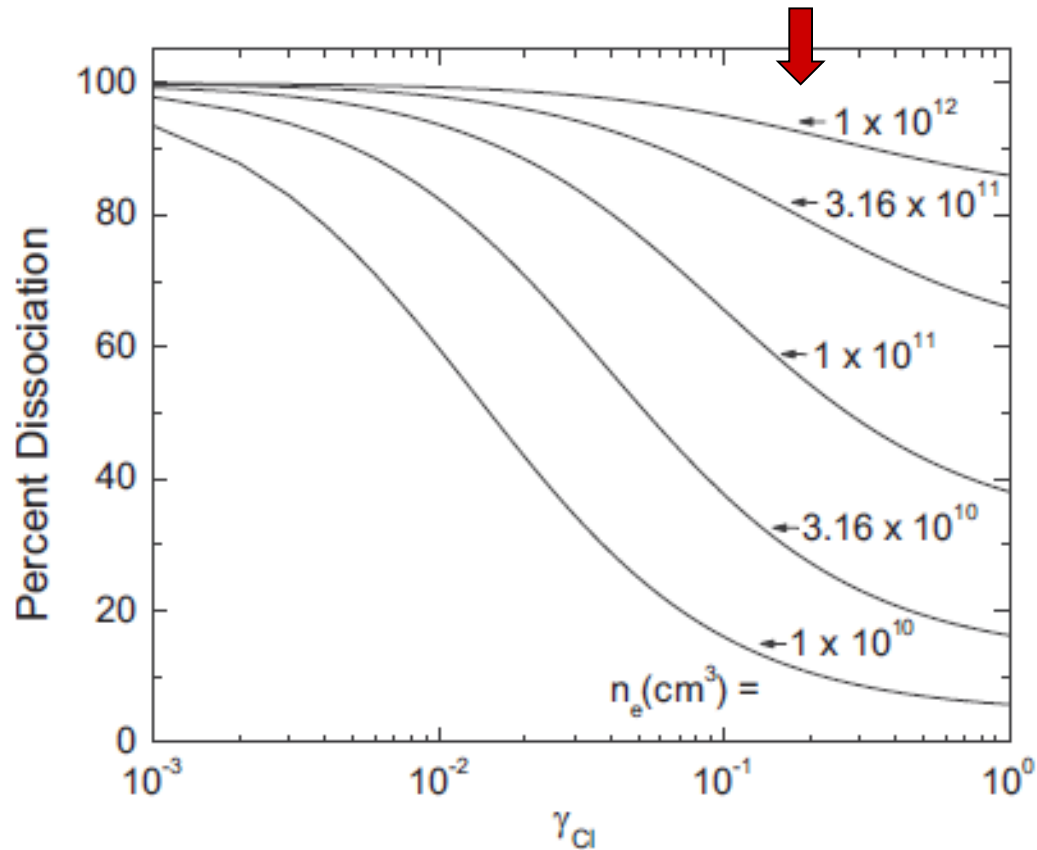
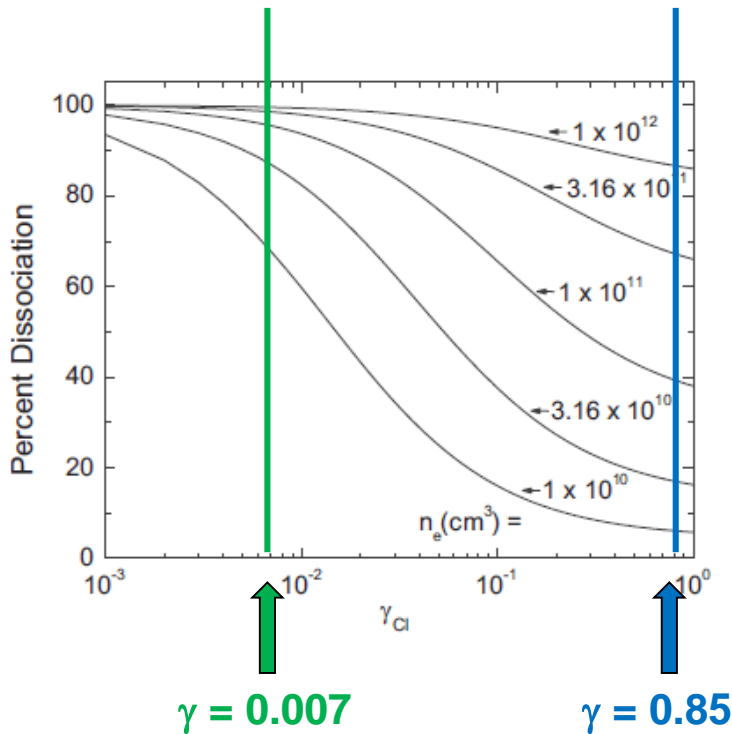


FIG. 1. Global model predictions of percent dissociation of  $\text{Cl}_2$  in a chlorine plasma as a function of the Cl recombination coefficient for different electron densities.  $R=L=20$  cm, pressure=10 mTorr, gas temperature,  $T_g=300$  K, and  $T_e=3$  eV.

# Values of gamma for Cl on stainless steel (literature)

V.M. Donnelly, J. Guha, L. Stafford, JVSTA. 010801(2011)



**Kota et al** : value derived from a high vacuum beam experiment :  $\gamma = 0.85$

**Richards and Sawin** : value derived from plasma measurements :  $\gamma = 0.007$

$\gamma$  dependent of :

- the nature of the material of the walls of the reactor
- the surface roughness (the specific area)
- the surface temperature
- the reactions at the wall : vacant sites created by ion bombardment.....

# How to determine $\gamma$ ?

best fit between predictions of the global model and experimental measurements in  $\text{Cl}_2$  ICP plasmas

⇒  $\gamma = 0.04$  and  $0.03$  in two different studies (Malyshev et al, Ullal et al)

## What happens at the walls ?

stainless steel walls are rapidly covered with silicon oxide or oxichloride ( $\text{SiOCl}$ ) due to etching of the dielectric quartz windows of the ICP

⇒ surface quite different from stainless steel

## General case :

$\gamma$  depends on the nature of the wall surface, which itself depends on the plasma process and species present in the plasma

⇒ need to measure  $\gamma$  in situ, in « real » deposition conditions

# Some examples of diagnostics used to determine in situ the value of $\gamma$

- **Absorption spectroscopy** → measurement of [X] and [X<sub>2</sub>]  
(G. Cunge et al for Cl and Cl<sub>2</sub>)
- **Fluorescence spectroscopy** → Radial profile of [X]  
(JP Booth et al, in CF<sub>4</sub> based plasmas)
- **Spinning wall method**  
(L. Stafford and V. Donnelly, in Cl<sub>2</sub> plasmas)
- **Time resolved optical emission spectroscopy in pulsed plasma**  
(A. Bouchoule and P. Ranson, A. Rousseau et al in H<sub>2</sub> plasmas, A. Granier et al in O<sub>2</sub> and O<sub>2</sub>/organometallic plasmas )

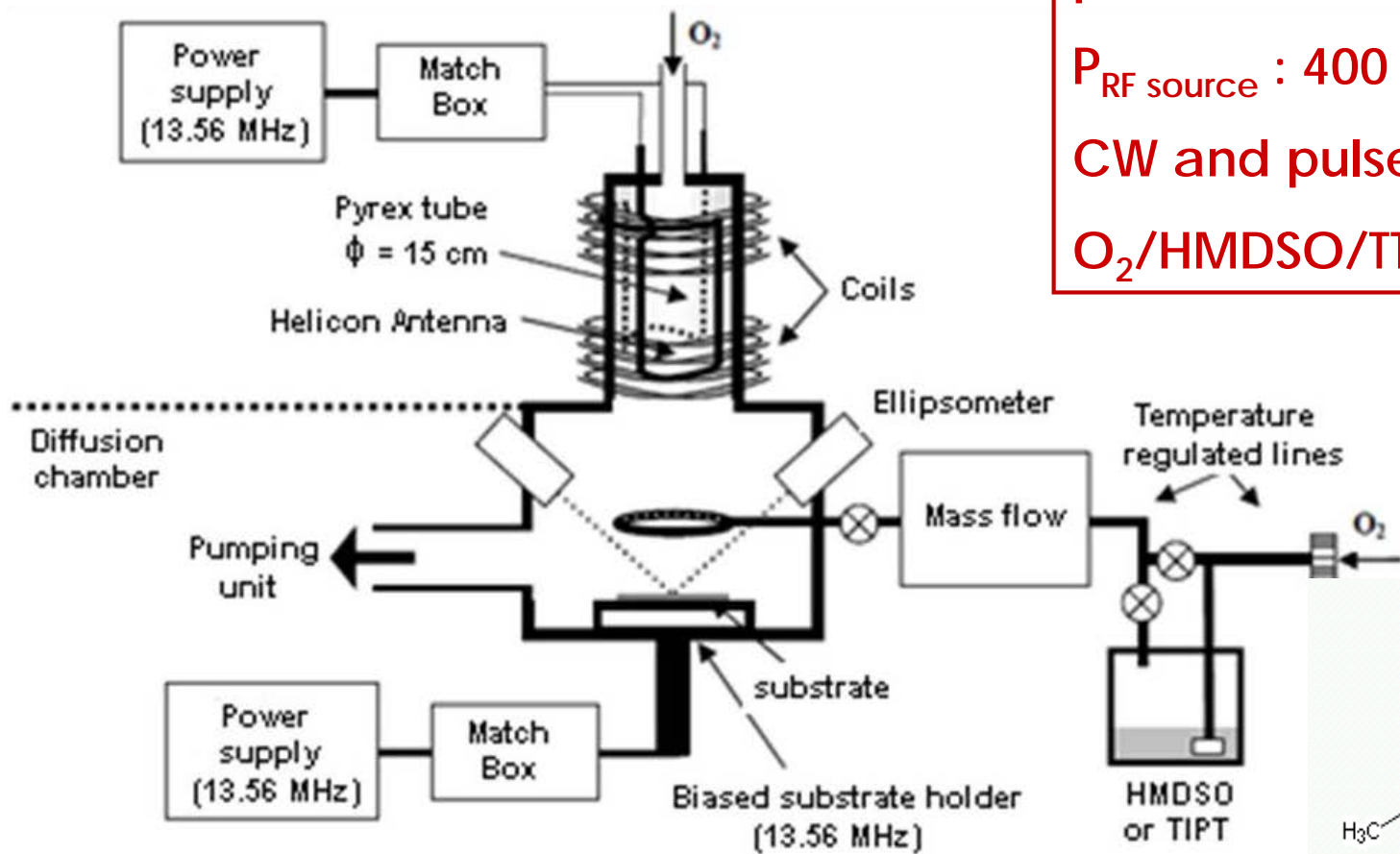
V.M. Donnelly, J. Guha, L. Stafford, JVSTA. 010801(2011)



# Exemple of determination of $\gamma$ in a low pressure ICP plasma by time resolved emission spectroscopy (TROES) in pulsed $O_2$ plasmas

# Low pressure RF ICP PECVD reactor

O<sub>2</sub>/Ar 95:5



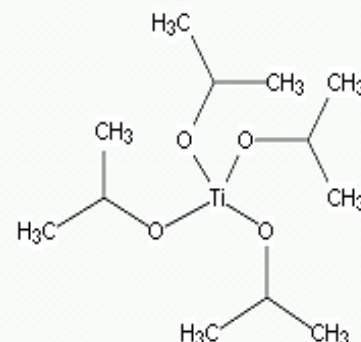
**p : 3 mTorr**

**P<sub>RF source</sub> : 400 W**

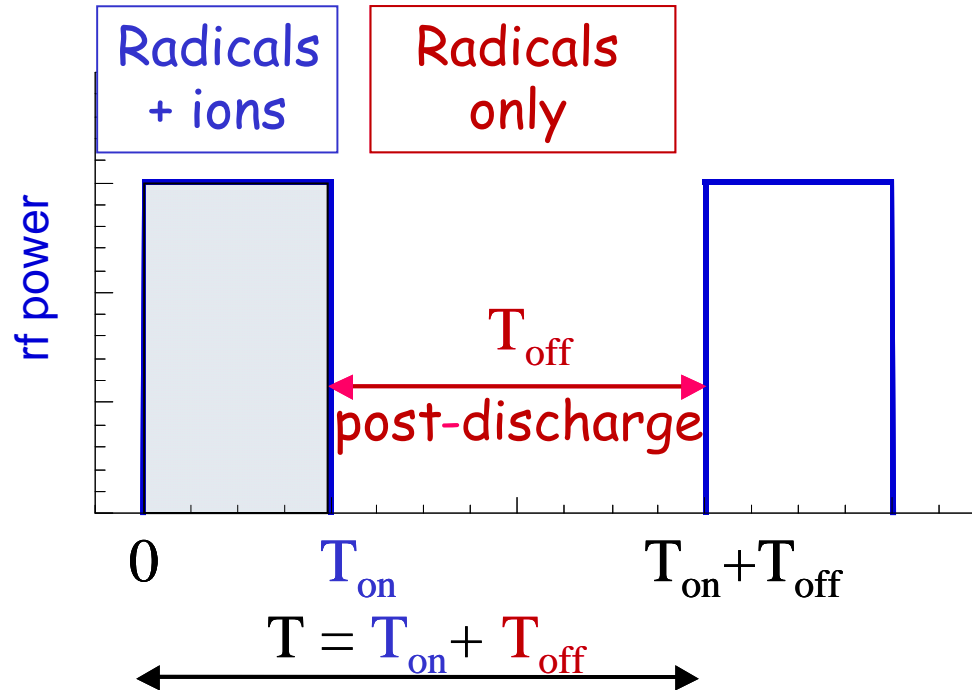
**CW and pulsed mode**

**O<sub>2</sub>/HMDSO/TTIP plasmas**

**TTIP**



# Characteristic times in pulsed plasmas

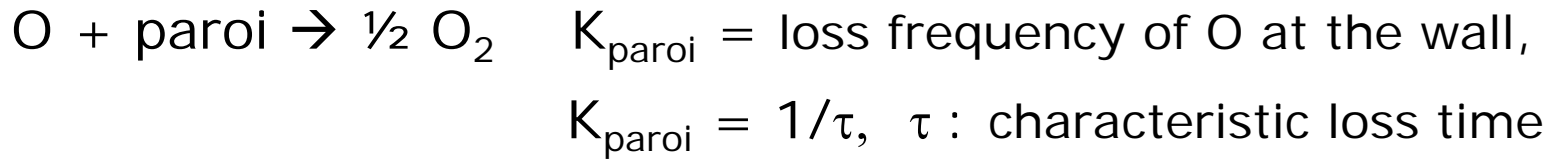


$$\text{DC} = \frac{T_{\text{on}}}{T_{\text{on}} + T_{\text{off}}}$$

- Characteristic times in low pressure plasmas ( $1 \text{ mTorr} < p < 1 \text{ Torr}$ ):
- Electron temperature rise and decay :  $10\text{-}50 \mu\text{s}$
- Radical lifetime :  $1\text{-}10 \text{ ms}$
- Residence time of stable species :  $100\text{-}1000 \text{ ms}$

**1 Hz - 100 kHz**

## O atom kinetics (O in fundamental state)



### Plasma created in continuous mode (CW) :

$$[\text{O}] = 2 K_{\text{diss}} \times n_e \times [\text{O}_2] / K_{\text{paroi}} = t \times K_{\text{diss}} \times n_e \times [\text{O}_2]$$

### Plasma created in pulsed mode :

$$[\text{X}]_{\text{ON}}(t) = [\text{X}]_{\text{max}} \cdot \left[ 1 - \exp\left(-\frac{t}{\tau_{\text{on}}}\right) \right]$$

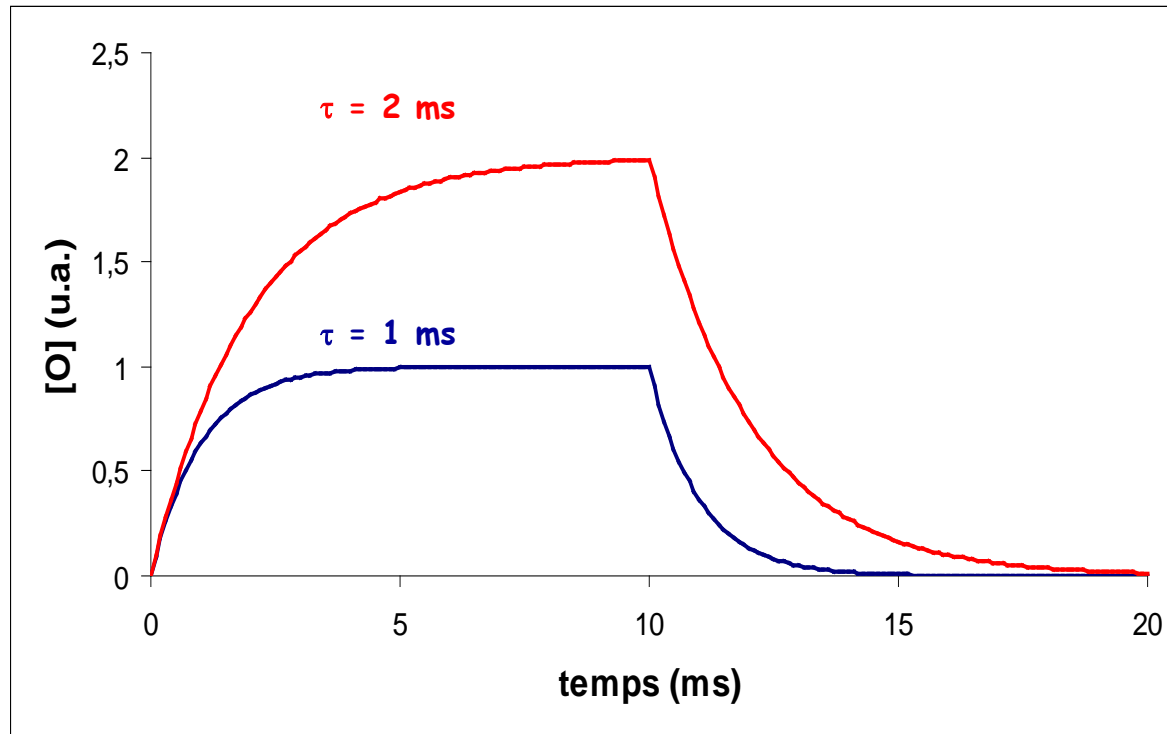
$$[\text{X}]_{\text{off}}(t) = [\text{X}]_{\text{max}} \cdot \exp\left(-\frac{t}{\tau_{\text{off}}}\right)$$

# Kinetics of X atoms in fundamental state

if  $\tau$  is doubled (i.e.  $K_{\text{paroi}}$  divided by 2)

→ creation of X 2 times more slowly

→ density of X at the stationary state multiplied by 2



# Determination of [O] from $I_O/I_{Ar}$

O<sub>2</sub> (+ 5% Ar) - O line at 844.6 nm and Ar line at 750.4 nm

hypotheses :

- maxwellian EEDF,  $T_e$  measured in O<sub>2</sub> plasma ( $T_e = 3.5$  eV)
- $O + e \rightarrow O^* + e$  et  $O_2 + e \rightarrow O^* + O + e$  (Walkup et al, 1986)
- $Ar + e \rightarrow Ar^* + e$  (Lawrence, 1976)
- O\* and Ar\* lost by radiative deexcitation
- $C(\lambda)$  response of the optical system (measured)

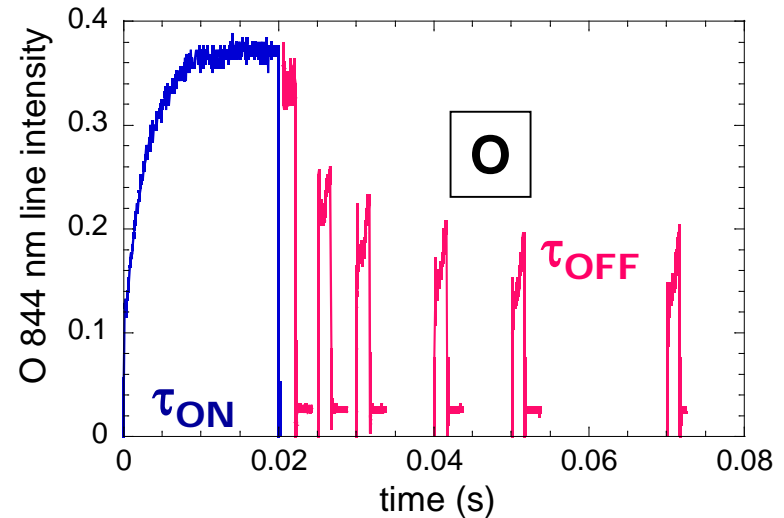
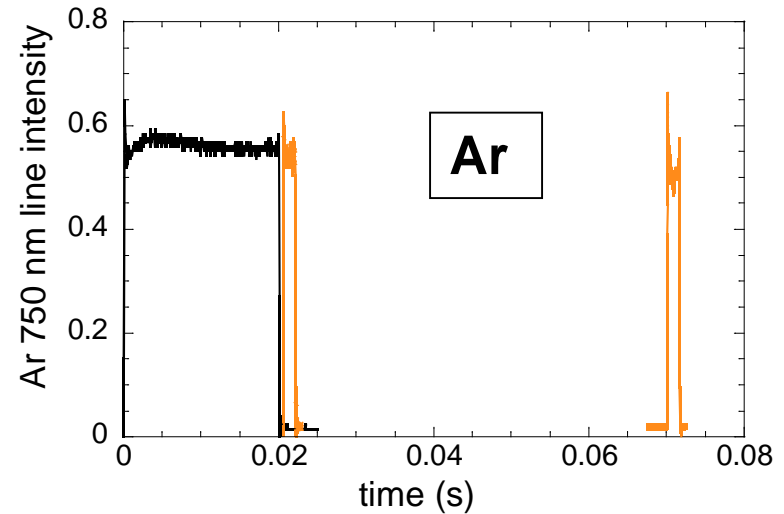
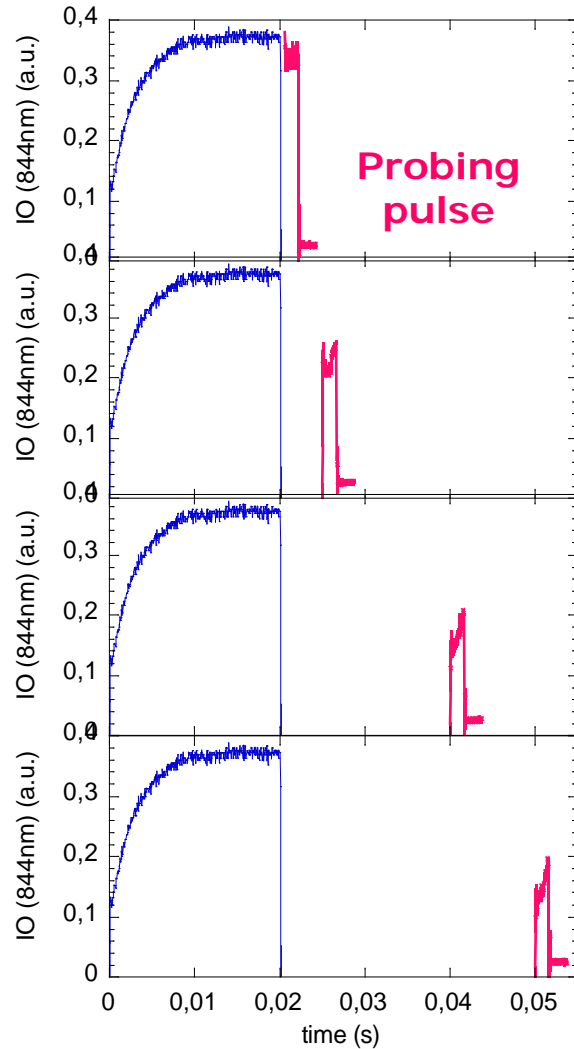


$$\frac{I_O}{I_{Ar}} = \frac{C(\lambda_O)}{C(\lambda_{Ar})} \times \frac{\nu_O}{\nu_{Ar}} \times \frac{k_{exc}^O [O] + k_{diss}^{O_2} [O_2]}{k_{exc}^{Ar} [Ar]}$$

A. Granier et al , J. Appl. Phys 75, 101 (1994)

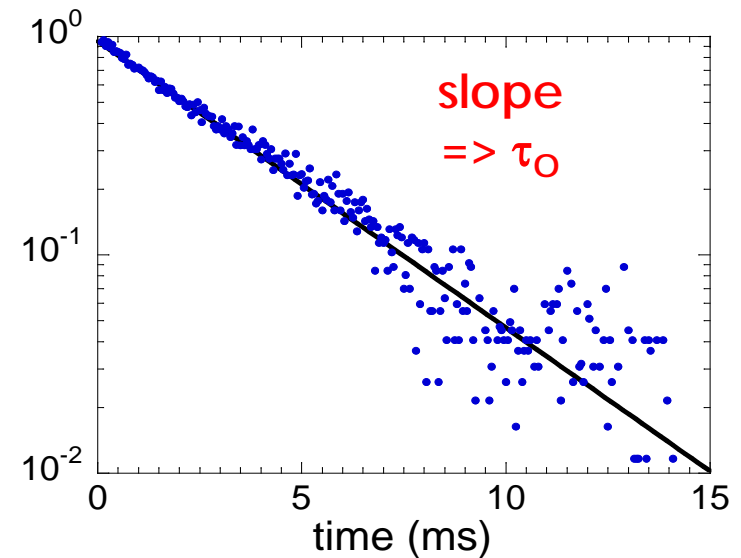
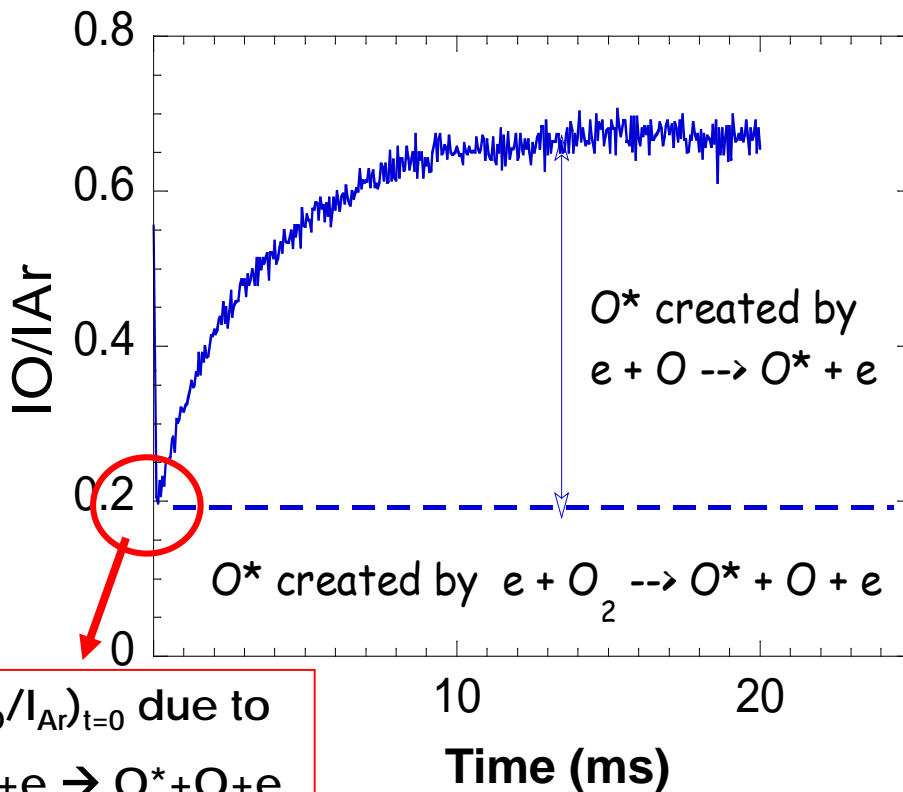
A. Granier et al, Plasma Sources Sci. Technol. 97, 147 (1997)

# Time Resolved Optical Emission Spectroscopy



A. Bouchoule, P. Ranson - J. Vac. Sci. Technol. A 9(2) (1991) 317.

# Determination of $\gamma_O$ from $I_{O\ 844nm}/I_{Ar\ 750\ nm}$ ( $T_{on}$ )



$$\tau_{on} = \tau_{diff} + \tau_{loss}$$

$$\tau_{on} = \frac{\Lambda_0^2}{D} + \frac{V}{A} \cdot \frac{2(2 - \gamma_{on})}{\bar{v} \cdot \gamma_{on}}$$

Chantry et al J. Appl. Phys. 62, 1141 (1987)

→  $\gamma_O$

$(I_O/I_{Ar})_{t=0}$  due to  
 $O_2 + e \rightarrow O^* + O + e$

$$\frac{I_O}{I_{Ar}} = \frac{C(\lambda_O)}{C(\lambda_{Ar})} \times \frac{v_O}{v_{Ar}} \times \frac{k_{exc}^O [O] + k_{diss}^{O_2} [O_2]}{k_{exc}^{Ar} [Ar]}$$

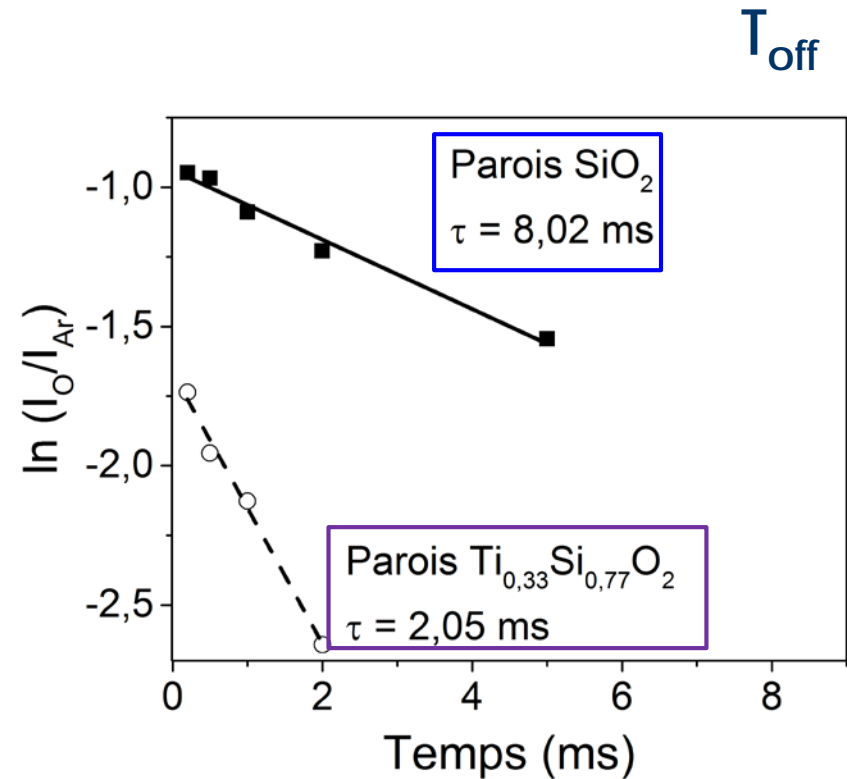
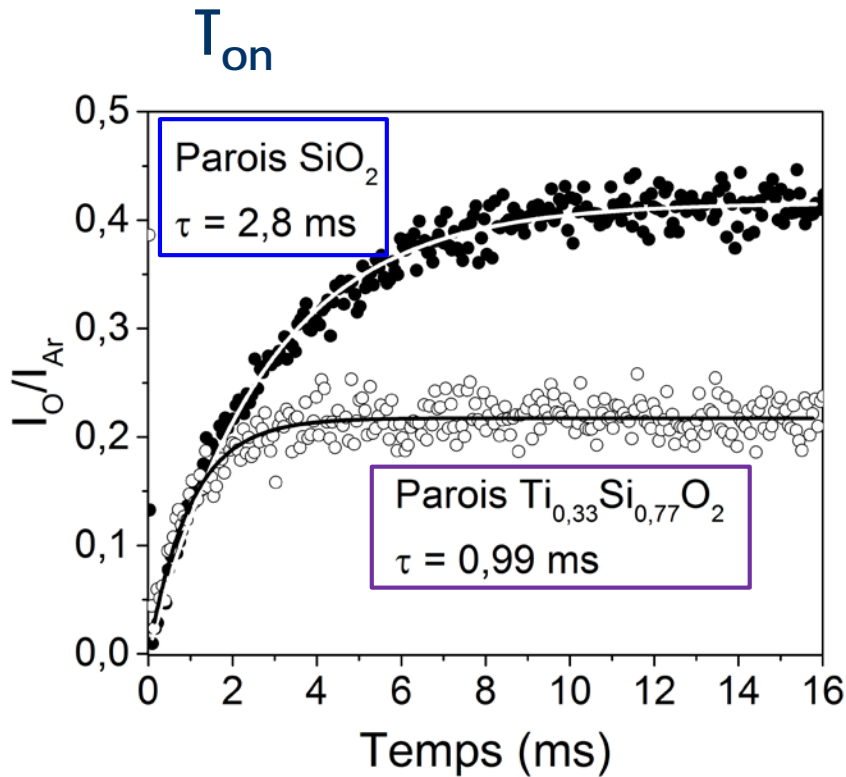


Oxygen plasma created under the same conditions (pressure, power) with the walls covered with different materials:

- $\text{SiO}_2$  (deposited in  $\text{O}_2$ /HMDSO plasmas)
- $\text{TiO}_2$  (deposited in  $\text{O}_2$ /TTIP plasmas)
- $\text{Ti}_x\text{Si}_{1-x}\text{O}_2$  (deposited) in  $\text{O}_2$ /HMDSO/TTIP plasmas

# Influence of the wall nature on O kinetics in O<sub>2</sub> plasmas (T<sub>on</sub>)

5 Hz, T<sub>on</sub> = 20 ms, T<sub>off</sub> = 180 ms



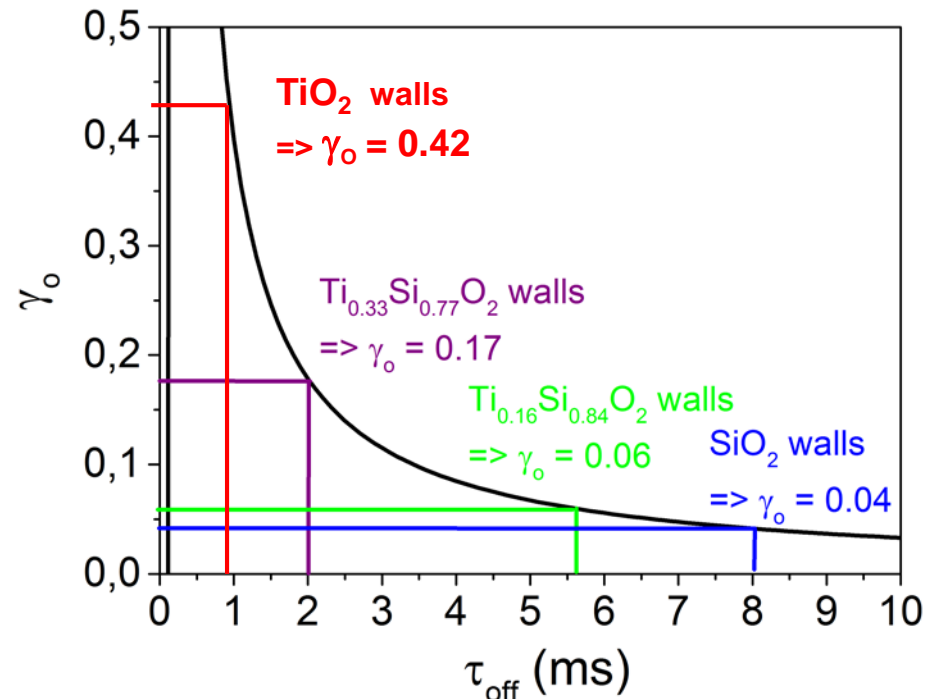
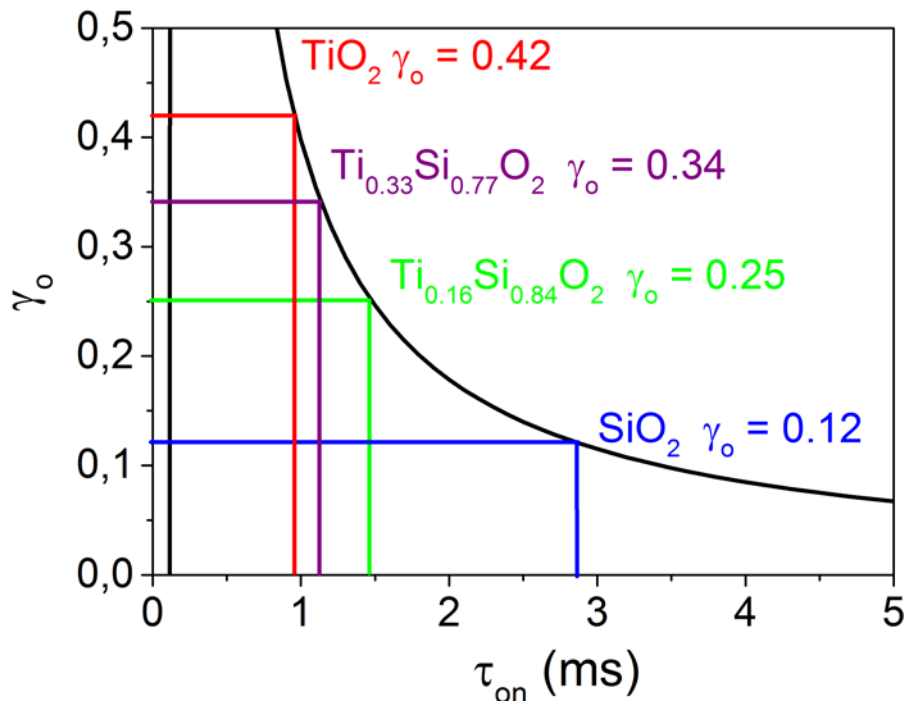
Thèse A. Bousquet (2005), Thèse de S. Elisabeth (2015)

A. Bousquet, A. Goulet, A. Granier dans "Plasmas Froids. Interactions Plasma-Surface : Modèles, Diagnostics et Procédés». Collection : Intégrations, Vol 7 – Ed. G Lelièvre, MRCT CNRS 2011 ISBN 978-2-918701-05-7 (2011), pp 211-234.

# Influence of the wall nature on $\gamma_o$ in $O_2$ plasmas ( $T_{on}$ & $T_{off}$ )

Chantry et al *J. Appl. Phys.* 62, 1141 (1987)

$$\tau = \tau_{diff} + \tau_{loss} = \frac{\Lambda_0^2}{D} + \frac{V}{A} \cdot \frac{2(2-\gamma)}{\bar{v} \cdot \gamma}$$



Thèse de S. Elisabeth (2015)

A. Bousquet, G. Cartry, A. Granier, *Plasma Sources Science and Technology* 16, 597-605 (2007)

# Influence of the wall nature on $\gamma_O$ in $O_2$ plasmas ( $T_{off}$ )

wall material	$\gamma_O (T_{on})$	$\gamma_O (T_{off})$
$TiO_2$	$0.42 \pm 0.2$	$0.42 \pm 0.2$
$Ti_{0.33}Si_{0.67}O_2$	$0.34 \pm 0.1$	$0.17 \pm 0.05$
$Ti_{0.16}Si_{0.84}O_2$	$0.25 \pm 0.1$	-
$SiO_2$	$0.12 \pm 0.02$	$0.04 \pm 0.01$

- $\gamma_O (T_{on})$  decreases as Si added to  $TiO_2$ , from 0.42 on  $TiO_2$  to 0.12 on  $SiO_2$
- $TiO_2$  : same (high) value of  $\gamma_O$  during  $T_{on}$  and  $T_{off}$
- $SiO_2$  :  $\gamma_O$  reduced by a factor 3 in the post-discharge

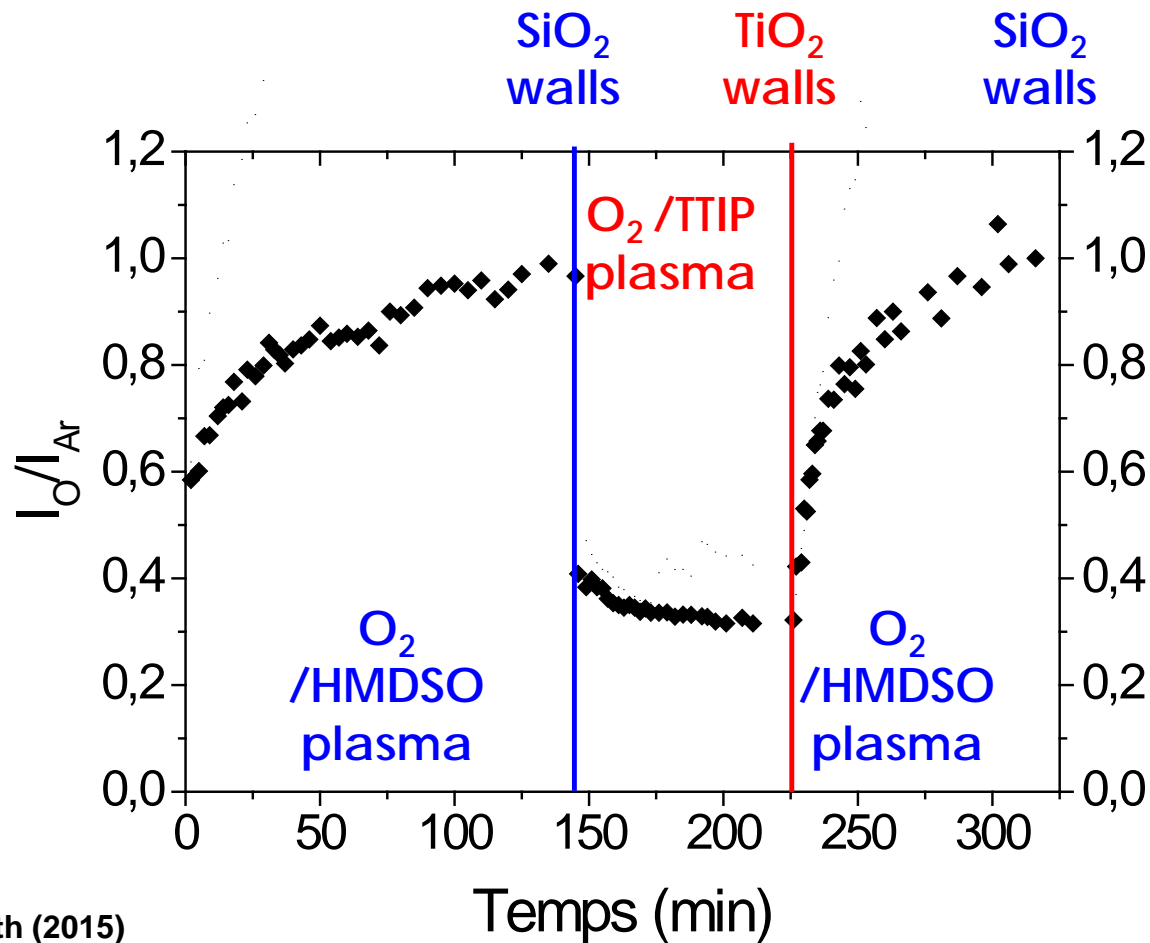
consistent with the literature and attributed to ion induced recombination and absence of ion bombardment during the post-discharge

Thèse A. Bousquet (2005), Thèse de S. Elisabeth (2015)

# O atom kinetics studied by TROES in pulsed O<sub>2</sub>/TIPT and O<sub>2</sub>/HMDSO plasmas

# In situ demonstration of the strong effect of the nature of the oxide covering the wall by OES : O atoms

$I_{O}/I_{Ar}$  monitored during successive cycles of  $O_2$ /HMDSO and  $O_2$ /TTIP plasmas



Thèse de S. Elisabeth (2015)

# O atom kinetics in O<sub>2</sub>/TTIP and O<sub>2</sub>/HMDSO plasmas

PECVD plasma	$\gamma_O (T_{on})$	O <sub>2</sub> plasma on	$\gamma_O (T_{on})$
O <sub>2</sub> /TTIP 98.5 : 1.5	0.42 ± 0.2	TiO <sub>2</sub>	0.42 ± 0.2
O <sub>2</sub> /HMDSO 99.5 : 0.5	0.065 ± 0.01	SiO <sub>2</sub>	0.12 ± 0.02

✦ TiO<sub>2</sub> walls : same value of  $\gamma_O$  in O<sub>2</sub>/TTIP plasmas and O<sub>2</sub> plasma

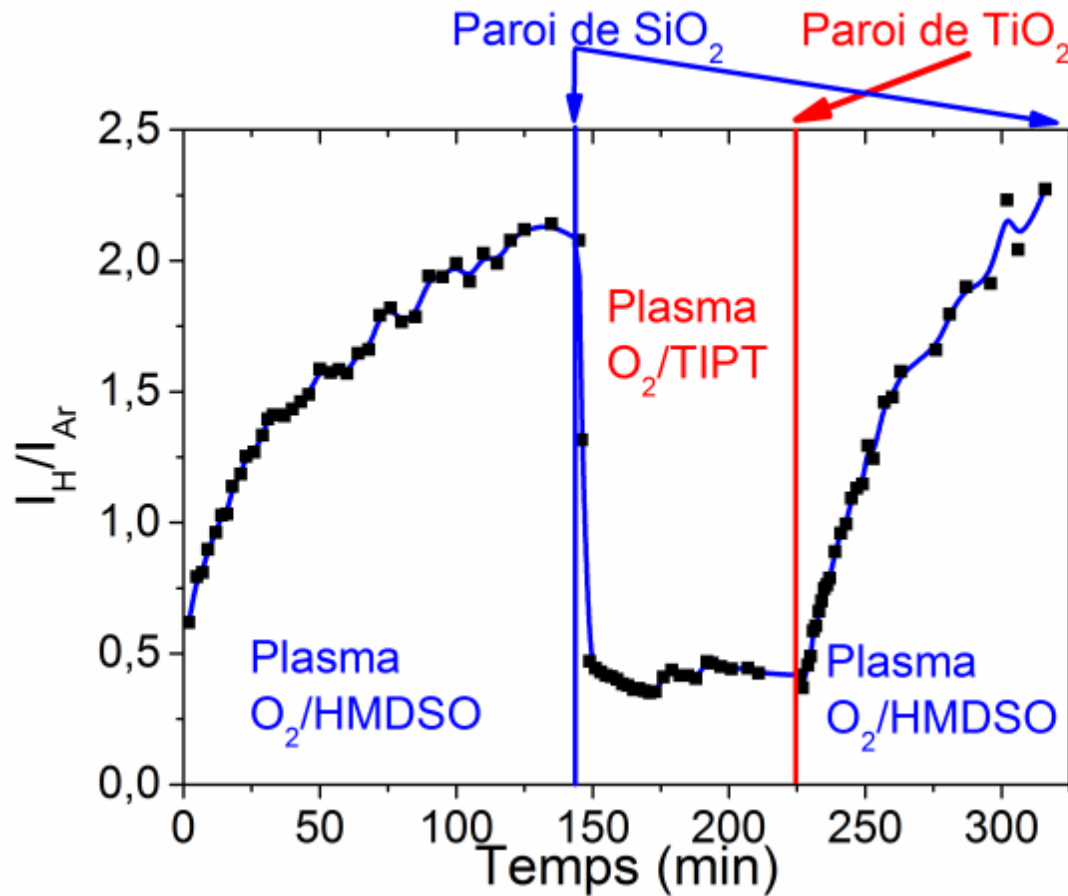
✦ SiO<sub>2</sub> walls :  $\gamma_O$  in O<sub>2</sub>/TTIP plasmas = 0.5 x  $\gamma_O$  in O<sub>2</sub> plasma

attributed to OH reactivity on reactive sites which limits O recombination in O<sub>2</sub>/HMDSO plasma (consistent with results obtained in O<sub>2</sub>/HMDSO 85:15 plasmas, where the effect was stronger due to higher OH contents)

**Does it mean that O is more reactive than OH on TiO<sub>2</sub> surfaces while OH more reactive than O on SiO<sub>2</sub> surfaces ?**

# In situ demonstration of the strong effect of the nature of the oxide covering the wall by OES : H atoms

$I_H/I_{Ar}$  monitored during successive cycles of  $O_2$ /HMDSO and  $O_2$ /TTIP plasmas



Thèse de S. Elisabeth (2015)



# H atom kinetics in O<sub>2</sub>/TTIP and O<sub>2</sub>/HMDSO plasmas

PECVD plasma	$\gamma_H (T_{on})$
O <sub>2</sub> /TTIP 98.5 : 1.5	0.065 ± 0.2
O <sub>2</sub> /HMDSO 99.5 : 0.5	0.01 ± 0.002

$\gamma_H$  in O<sub>2</sub>/TTIP plasmas (TiO<sub>2</sub> walls) = 6 x  $\gamma_H$  in O<sub>2</sub> /HMDSO plasma (SiO<sub>2</sub> walls)

# Conclusion

Real difficulty to predict values of recombination frequencies

Atom densities can vary on very large scales according to the nature of the walls for the same discharge parameters (power, pressure...), typically from 1% to several ten percent

Pulsed plasma : for the same « walls »  $\gamma$  values most often different in discharge and post-discharge times